

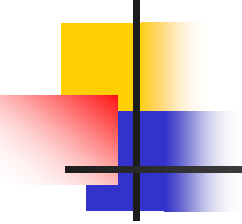


AN INTRODUCTION TO QUANTUM OPTICS...

...the light as you've never seen
before...

Optics and Photonics
LCI – Spring 2008
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*"If things of sight such heavens be, what
heavens are those we cannot see?"*

Andrew Marvell

INTRODUCTION

- First steps:
- Quantum by Plank;
- Photon by Einstein;

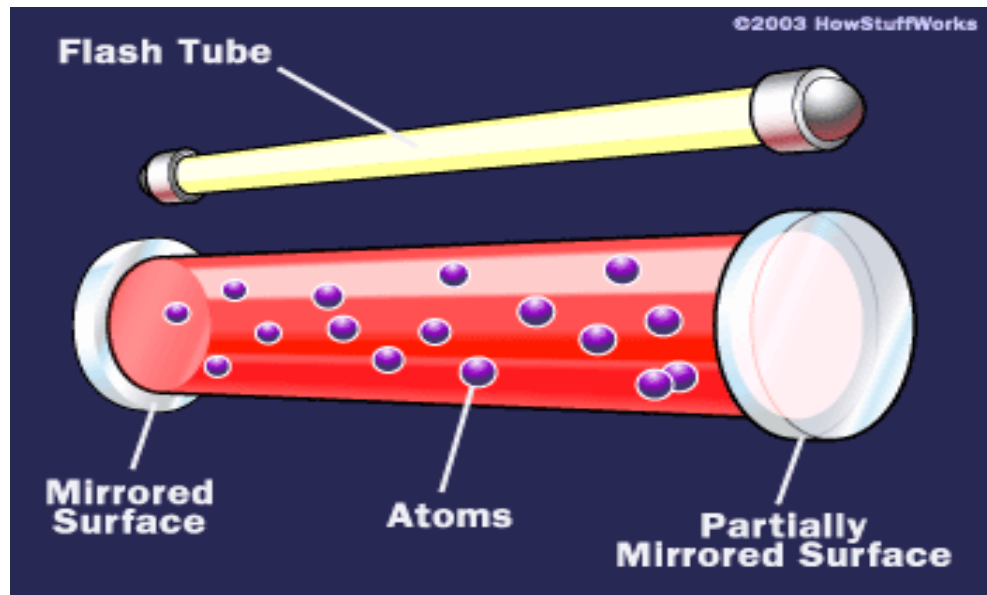
Newton X Huygens

- Feynman "ideal" experiment;
- And then, for the next 55 years...



INTRODUCTION

- And the **L**ight **A**mplification by **S**timulated **E**mission of **R**adiation arrives:





DEFINITION

- Quantum Optics:

“Quantum optics is a field in quantum physics, dealing with the application of quantum mechanics to phenomena involving light and its interactions with matter.”



QUANTUM OPTICS OPERATORS

- Light is described in terms of field operators for creation and annihilation of photons:

$$a^\dagger = \frac{1}{\sqrt{2}}(q - ip) \quad \textit{creation operator or raising operator}$$

$$a = \frac{1}{\sqrt{2}}(q + ip) \quad \textit{annihilation operator or lowering operator}$$

$$H = \hbar\omega \left(a^\dagger a + \frac{1}{2} \right)$$



QUANTUM FIELDS

- The usual description of the classical fields are given by quantum versions:

$$\hat{E}_x(z, t) = \mathcal{E}_0(a + a^\dagger)\sin(kz)$$

$$\hat{B}_y(z, t) = \mathcal{B}_0\frac{1}{i}(a - a^\dagger)\cos(kz)$$

where \mathcal{E}_0 and \mathcal{B}_0 are the electric and magnetic field
“per photon”.



QUANTUM FIELDS

- The number state has well-defined energy but not well defined electric field:

$$\langle n | \hat{E}_x(z, t) | n \rangle = \mathcal{E}_0 \sin(kz) [\langle n | a | n \rangle + \langle n | a^\dagger | n \rangle] = 0$$

but the fluctuations of the electric field are not zero:

$$\Delta E_x = \sqrt{2\mathcal{E}_0} \sin(kz) \left(n + \frac{1}{2} \right)^{1/2}$$



THE QUANTUM PHASE

- The single mode electric field is:

$$\hat{E}_x(r, t) = i \left(\frac{\hbar\omega}{2\varepsilon_0 V} \right)^{1/2} e_x (a e^{ik \cdot r - i\omega t} + a^\dagger e^{-ik \cdot r + i\omega t})$$

the classical theory:

$$\hat{E}_x(r, t) = \frac{1}{2} E_0 e_x (e^{ik \cdot r - i\omega t + \phi} + e^{-ik \cdot r - i\omega t + \phi})$$

May a and a^\dagger have information about the phase?



THE QUANTUM PHASE

- The Susskind-Glogower (SG) operators:

$$\hat{E} = (aa^\dagger)^{-1/2}a \quad \text{and} \quad \hat{E}^\dagger = a^\dagger(aa^\dagger)^{-1/2}$$

are the quantum phase factors with eigenstates:

$$\hat{E}|\phi\rangle = e^{i\phi}|\phi\rangle$$

$$|\phi\rangle = \sum_{n=0}^{\infty} e^{in\phi}|n\rangle$$



COHERENT STATE – THE MOST CLASSICAL QUANTUM STATE

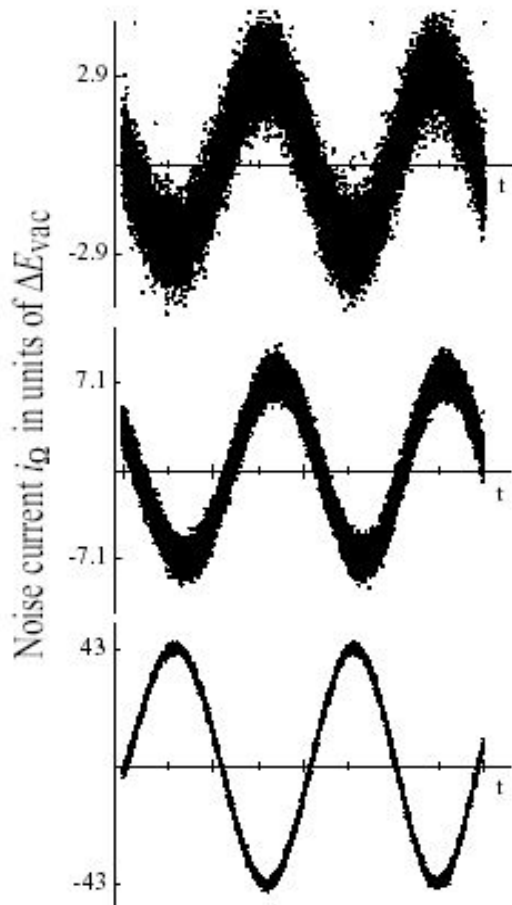
- The expected value of the “classical” electric field is not zero – single states are not the “real” description. Why not define:

$$|\alpha\rangle = \sum_{n=0}^{\infty} C_n |n\rangle$$

which give us the classical expected value of the electric field. And...

$$\langle (\Delta X_1)^2 \rangle_{\alpha} = \frac{1}{4} = \langle (\Delta X_2)^2 \rangle_{\alpha}$$

COHERENT STATE – THE MOST CLASSICAL QUANTUM STATE

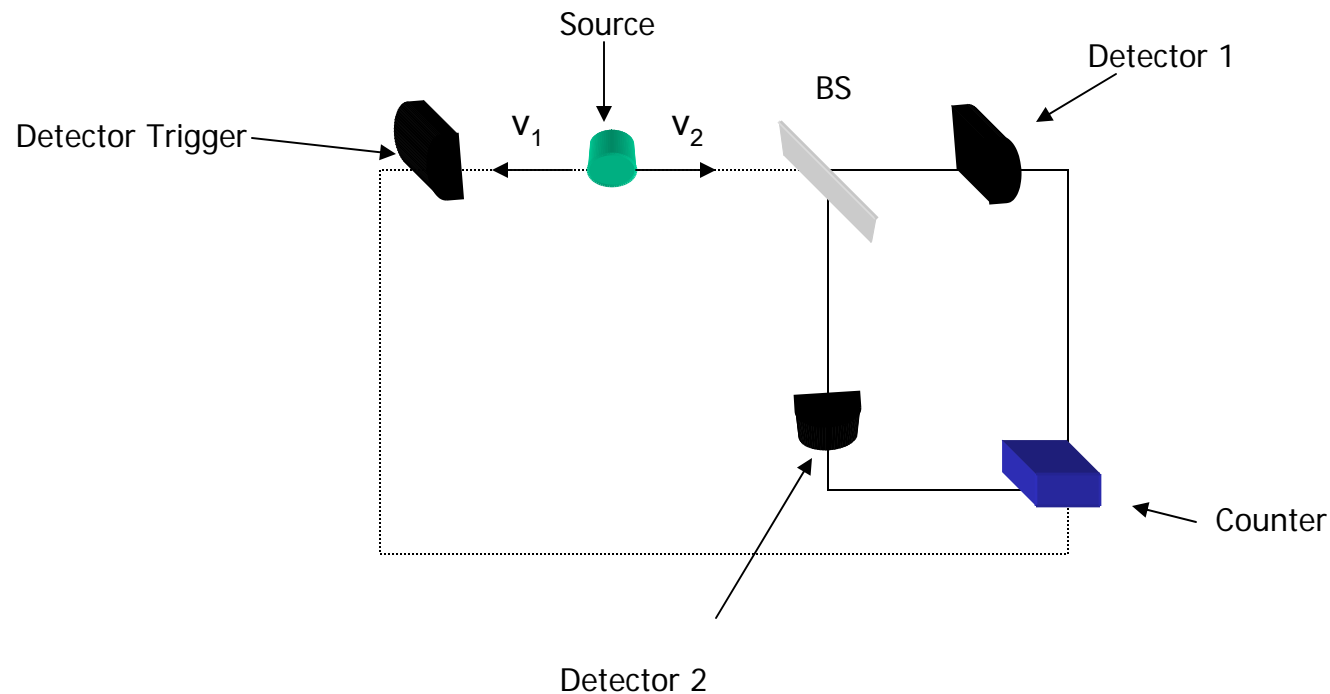


The operators “phase” and “amplitude” do not commute and the lower limit is the coherent state:

LASER

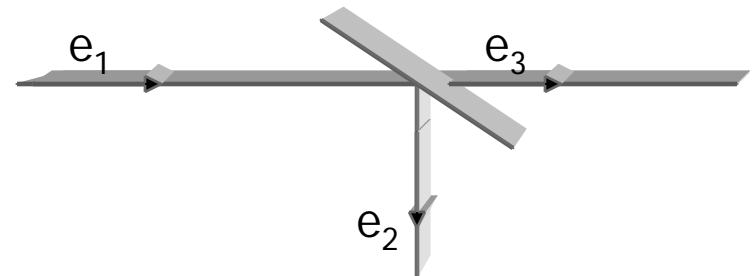
A LITTLE BIT OPTICS – BEAM SPLITTERS AND INTERFEROMETERS

Grangier experiment:



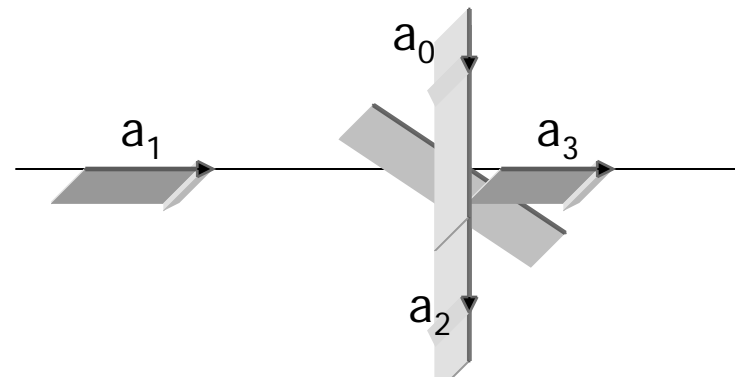
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Classical beam splitter

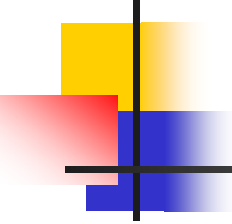


$$e_2 = re_1, \quad e_3 = te_1 \quad \text{and} \quad |r|^2 + |t|^2 = 1$$

Quantum beam splitter



$$a_2 = ra_1 + t'a_0, \quad a_3 = ta_1 + r'a_0$$



A LITTLE BIT OPTICS – BEAM SPLITTERS AND INTERFEROMETERS

If we have a 50:50 beam splitter:

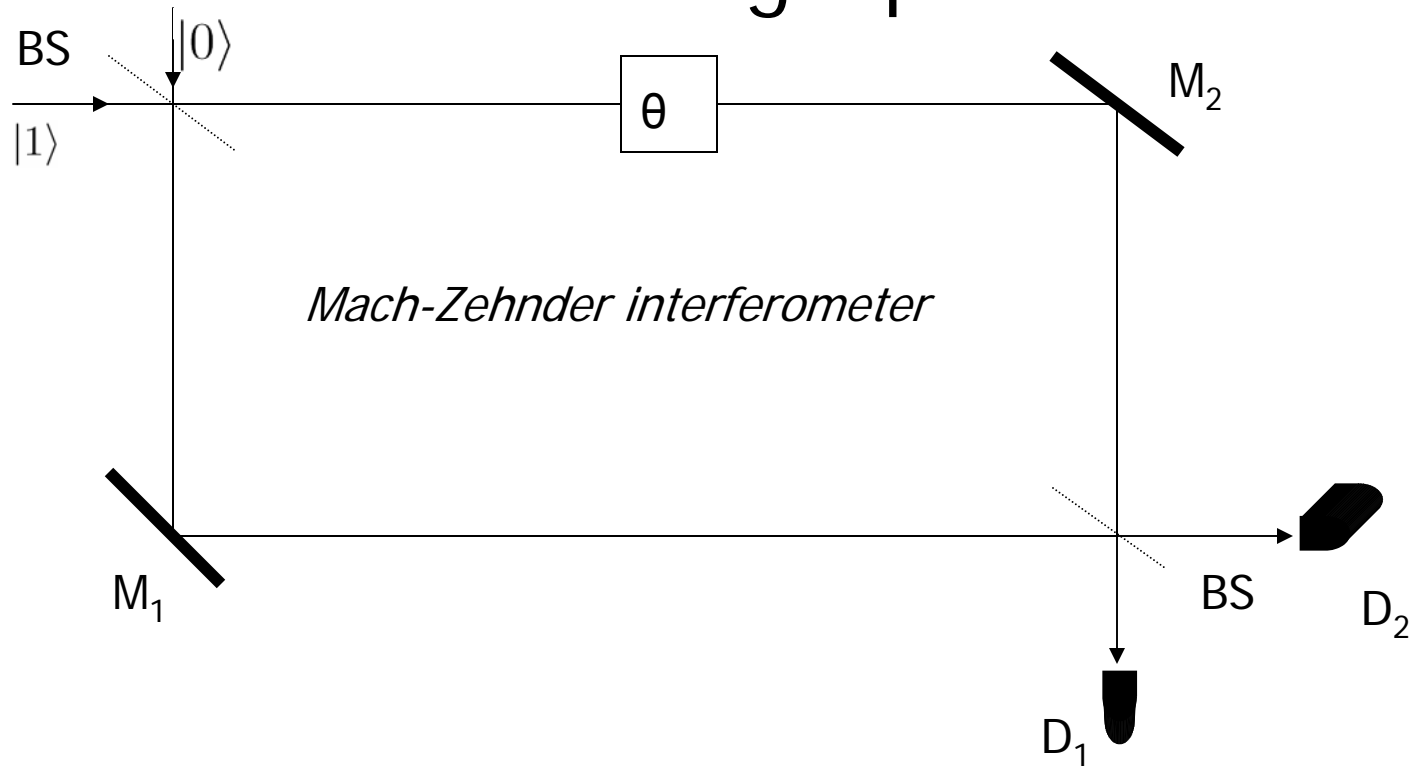
$$a_2 = \frac{1}{\sqrt{2}}(a_0 + ia_1), \quad a_3 = \frac{1}{\sqrt{2}}(ia_0 + a_1)$$

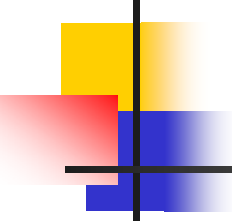
so, given an incoming state $|0\rangle_0|1\rangle_1$

$$|0\rangle_0|1\rangle_1 \xrightarrow{\text{BS}} \frac{1}{\sqrt{2}}(i|1\rangle_2|0\rangle_3 + |0\rangle_2|1\rangle_3)$$

A LITTLE BIT OPTICS – BEAM SPLITTERS AND INTERFEROMETERS

■ Interference of single photons:





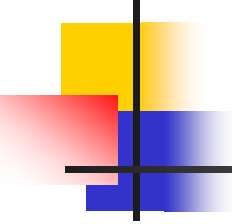
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- For an incoming state $|0\rangle_0|1\rangle_1$:

$$|0\rangle_0|1\rangle_1 \xrightarrow{\text{BS}} \frac{1}{\sqrt{2}}(i|1\rangle_2|0\rangle_3 + |0\rangle_2|1\rangle_3)$$

the phase shifter is represented by

$$|\phi\rangle = \sum_{n=0}^{\infty} e^{in\phi}|n\rangle$$



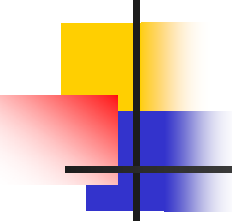
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SO...

$$\frac{1}{\sqrt{2}}(i|1\rangle_2|0\rangle_3 + |0\rangle_2|1\rangle_3) \xrightarrow{\text{BS}} \frac{1}{\sqrt{2}}(i|1\rangle_2|0\rangle_3 + e^{i\theta}|0\rangle_2|1\rangle_3)$$

and, after the next beam splitter

$$\frac{1}{\sqrt{2}}(i|1\rangle_2|0\rangle_3 + |0\rangle_2|1\rangle_3) \xrightarrow{\text{BS}} \frac{1}{\sqrt{2}} [(e^{i\theta} - 1)|0\rangle_2|1\rangle_3 + i(e^{i\theta} + 1)|1\rangle_2|0\rangle_3]$$



A LITTLE BIT OPTICS – BEAM SPLITTERS AND INTERFEROMETERS

then we have states with probability:

$$P_{01} = \frac{1}{2}(1 - \cos \theta)$$

$$P_{10} = \frac{1}{2}(1 + \cos \theta)$$

Nonlocality – detection without interaction!



AND QUANTUM OPTICS...

*...The squeezed states of light (quantum noise);
Optical trap; Quantum entanglement; Quantum
teleportation, Single photon experiments;
Interactions of single photon and matter (optical
cavity); Quantum networks; Quantum
communication; Bose Einstein condensates;
Quantum encryption and so on...*

OH ALICE... YOU'RE
THE ONE FOR ME

BUT BOB... IN A
QUANTUM WORLD
HOW CAN WE BE SURE?

ψ^+ or ψ^- ?



“The word classical means only one thing in science: it’s wrong!” — B.R. Frieden



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